

Effect of Laser Shock Peening (LSP) on the Phase Evolution, Residual Stress and Hardness of Hastelloy-X Superalloys

S. NATH*, P. SHUKLA, X-J. SHEN AND J. LAWRENCE

*School of Mechanical, Aerospace and Automotive Engineering, Faculty of Engineering,
Environment and Computing, Coventry University, Gulson Road, Coventry, CV1 2JH, UK*

Efforts have been made to understand the effect of laser shock peening (LSP) on the phase, residual stress and hardness of Hastelloy-X superalloys. A 10 J Nd:YAG laser was used for the LSP operation. Following LSP detailed microstructural and phase analysis along with residual stress and hardness studies were undertaken. A parametric window was first established to explore the relationship between LSP process parameters and the respective surface and bulk properties. The effects of an absorptive layer on the properties of the modified layer were also investigated. Qualitative and quantitative information on dislocation density was obtained using X-ray diffraction (XRD) analysis and correlated with the processing parameters. Residual stress developed following LSP was measured using the XRD technique. An increase in the hardness of the Ni-based superalloys was observed. The residual stress on the surface of the laser shock peened Hastelloy-X superalloy showed a maximum compressive stress of 166 MPa. A detailed microstructure-property relationship was established to understand the mechanism of property enhancement. Further optimization of the LSP process to surface treat the Hastelloy-X superalloys will open up new avenues for the material's applicability, particularly in the aerospace sector.

Keywords: Nd:YAG laser, Hastelloy-X, superalloy, absorptive layer, laser shock peening (LSP), residual stress, Hugoniot elastic limit (HEL), hardness; microstrain, dislocation density

*Corresponding author: E-mail: ac5689@coventry.ac.uk, sendsubha@gmail.com

1 INTRODUCTION

Laser shock peening (LSP) is a process normally used to induce compressive residual stress on the surface of the components that are exposed to a cyclic loading in a normal or corrosive environments [1-4]. The compressive residual stress helps to delay the crack initiation and propagation rate [5, 6]. The LSP treatment, under optimised process parameters, strain hardens the treated zone which in turn improves the tribological properties of the metallic systems [4]. The LSP treatment has also been tried on to ceramic systems [7].

Ni-based superalloys are a class of high temperature alloys which shows exceptional strength and resistance in terms of fatigue, creep, and corrosion [8]. Their superior performance at high temperature, allow them to be applicable for engine components in aerospace and automotive industry [8]; however, the demands for increasing the engine efficiency have pushed the engine to its extreme operating conditions. The changing operating conditions have increased the thermal and mechanical loads on the system, thereby, lowering their service life. Thus, strengthening of these materials is of great importance as it increases the efficiency and decreases the repair cost of the engines. The term strengthening means to improve the materials mechanical properties (hardness, fatigue strength, creep strength, etc.). It is very well known that the strength of materials can be improved by introducing foreign atoms in the matrix as a solid solution, *via* the effect of work hardening, by the formation of precipitates as a result of alloying, and decreasing the grain size of the matrix [9]. On the other hand, fatigue strength in materials can be improved by introducing compressive residual stress (CRS) on the surface of the component [1, 10]. The introduction of CRS on the surface of these components are normally achieved by shot peening and LSP [1, 11].

The aim of this investigation was to study the effect of LSP on the properties of the Hastelloy-X superalloys and attempts to further understand and verify the strengthening mechanisms as result of LSP surface treatment. In particular, the effect of LSP on the phase distribution, microstrain and residual stress evolution was studied using X-ray diffraction (XRD) to understand the contribution of LSP process parameters. Studies were also extended to elucidate the effect of absorptive layer on the development of residual stress on the surface of the superalloys. Finally, dislocation densities were measured to understand the work hardening behaviour of the superalloys followed by a study on the change in the microhardness.

2 EXPERIMENTAL DETAILS

2.1 Materials specifications and sample preparation

The LSP was carried out on rectangular Hastelloy-X superalloys samples of $30 \times 10 \times 10 \text{ mm}^3$ comprising of the composition shown in Table 1. Four

TABLE 1
Composition (wt.%) of the Hastelloy-X used for LSP.

Ni	Cr	Fe	C	Mo	Si	W
Bal.	20	19.8	7.6	7	0.3	0.5

superalloy coupons were used for the LSP. The superalloy samples were fully annealed at 1175° for 1 hour to relieve any machining stress. The superalloy samples were then polished using 600 µm grit size SiC papers to clean the surface oxides and residue. The superalloy samples were then ultrasonically cleaned using both acetone and isopropyl alcohol to remove any trapped carbide particles and any other contaminants.

2.2 Laser shock peening (LSP) apparatus and procedures

A 10 J, pulsed Q-switched Nd:YAG laser system (LPY10J; Litron, Ltd.) was used for the experimental study herein. The laser emitted a Gaussian beam at a wavelength of 1064 nm. The laser delivered 8 ns long pulse with a repetition rate of 10 Hz. The laser energy density ranged from 2 to 8 GW/cm² by keeping a constant input laser energy of 8 J, but altering the laser spot diameter. The beam divergence of the laser was 0.5 mrad with an M² value of 1.99. The combined laser beam characteristics and the applied process parameters exhibited a radiance density (laser beam brightness) ranging from 6.44 to 22.651.44 Jcm²/Sr¹µm, determined by the tried and tested methodology from the previous work [12-15], as detailed in Table 2.

Two samples were laser shock peened without absorptive coatings (LSPwc-1 and LSPwc-2) and the other two samples were laser shock peened with absorptive coatings (LSP-3 and LSP-4), as shown in Table 2. Two types of absorptive layers were used during LSP: one is an Al tape and the other is a black vinyl tape. All the samples were treated with an overlap of 50%.

According to the principle of LSP, the laser irradiation on the surface of the sample forms an expanding plasma with high pressure shock waves driven into the material. The shock waves with pressure greater than the Hugoniot elastic limit (HEL) of the material causes the material to deform plastically. According to Fabbro *et al* [16] and Peyre *et al* [2], the peak plasma pressure, P , can be expressed as

$$P = 0.01 \sqrt{\frac{\alpha}{2\alpha + 3}} \sqrt{z} \sqrt{I_0} \quad (1)$$

where I_0 is the laser power density, α is the efficiency of the interaction and Z is the reduced shock impedance between the material and the confining medium.

TABLE 2

Laser shock peening parameters used in the present study.

Sample	Energy (J)	Spot Size (mm)	Peak Power Density (GW/cm ²)	Radiance Density (Brightness, J.cm ² /Sr.μm)	Absorptive Layer	Overlap (%)	Peak Pressure (GPa)	HEL (GPa)
LSPwc-1	8	4.0	7.96	22.65	No layer	50	2.88	1.11
LSPwc-2	8	7.5	2.26	6.44	No layer	50	1.53	1.11
LSP-3	8	7.5	2.26	6.44	Aluminium tape	50	1.53	1.11
LSP-4	8	7.5	2.26	6.44	Black vinyl tape	50	1.53	1.11

The HEL of any material is related to its dynamic yield strength, (σ_Y^{dyn}), by [17]

$$HEL = \frac{1 - \nu}{1 - 2\nu} \sigma_Y^{dyn} \quad (2)$$

where ν is the Poisson's ratio of the material.

The laser shock peened (laser shock peened) parameters were chosen to ensure that the peak plasma pressure exceeds the HEL of the material to plastically deform the material. The laser shock peened parameters along with peak plasma pressure which was calculated using Equation (1) and HEL which was calculated using Equation (2) are presented in Table 2.

2.3 Material characterisation techniques

2.3.1 Phase analysis

A detailed analysis of the phase evolution was carried out by X-ray diffraction (XRD, D8 Discover; Bruker Corporation) with a scanning speed of 0.05°/s and time *per* step of 0.1 seconds using the Cu K α radiation.

2.3.2 Residual stress measurement

Residual stress developed on the surface of the thin films was carefully measured by the XRD using a Bragg-Brentano diffractometer. The X-ray source (Cu K α radiation) was operated at an accelerating voltage of 40 kV and current of 25 mA. The XRD was operated at a scanning speed of 0.01°/s and a scanning time of 5 seconds *per* step. The selected Ψ values were 0, 5, 10, 15, 20, 25, 30, 35, 40 and 45°. For the calculation of residual stress the (331) peak of γ -Ni phase was considered. The $\sin^2 \Psi$ technique was employed to measure the residual stress values in both the longitudinal and transverse directions. The equation employed to measure residual stress in any given direction, ϕ , is given by

$$\sigma_{\phi} = m \frac{1}{d_0} \left(\frac{E}{1 - \nu} \right)_{hkl} \quad (3)$$

where m is the slope of the d versus $\sin^2\psi$ curve, d_0 is the stress-free lattice spacing ($d_0 \approx d\phi_0$, where $d\phi_0$ is the stress free interplanar lattice spacing), E is the Young's modulus and ν is the Poisson's ratio. The suffix, hkl , refers to a particular crystallographic plane.

2.3.3 Dislocation density measurement

The dislocation density evolution, in the untreated and laser shock peened samples, was calculated using Williamson and Smallman approach [18]. Williamson and Smallman approach assumes the size and strain are related to dislocation density which is true when mechanical milling is used. However, in the present case as LSP was used to plastically deform the material, microstrain is the only parameter which is affected by plastic deformation and is related to dislocation density. The dislocation density due to strain in the material was calculated from

$$\rho = \frac{k\varepsilon^2}{b^2} \quad (4)$$

where $k = 16.1$ for a face-centred cubic (FCC) crystal; ε is the microstrain; and b is the burger vector which for an FCC crystal is $(a/2) \langle 110 \rangle$, where a is the lattice parameter. The lattice parameter in Hastelloy-X superalloy was measured to be 0.255 nm.

2.3.4 Microhardness measurement

The hardness of the untreated and LSP treated samples were measured by Vickers microhardness tester (DURASCAN-70; Struers, Inc.) with 50 mN load and a dwelling time of 10 seconds.

3 RESULTS AND DISCUSSION

3.1 Phase analysis

XRD phase scans on the untreated Hastelloy-X superalloy (Plot 1) and laser shock peened Hastelloy-X superalloy samples at 7.96 GW/cm² with no absorptive tape (Plot 2), 2.26 GW/cm² with no absorptive tape (Plot 3), 2.26 GW/cm² with Al as an absorptive tape (Plot 4), and 2.26 GW/cm² with black vinyl tape as an absorptive tape (Plot 5) are shown Figure 1. The XRD phase scans confirm the presence FCC Ni as the only phase (matrix) in both untreated and laser shock peened samples. The LSP treatment on the Hastel-

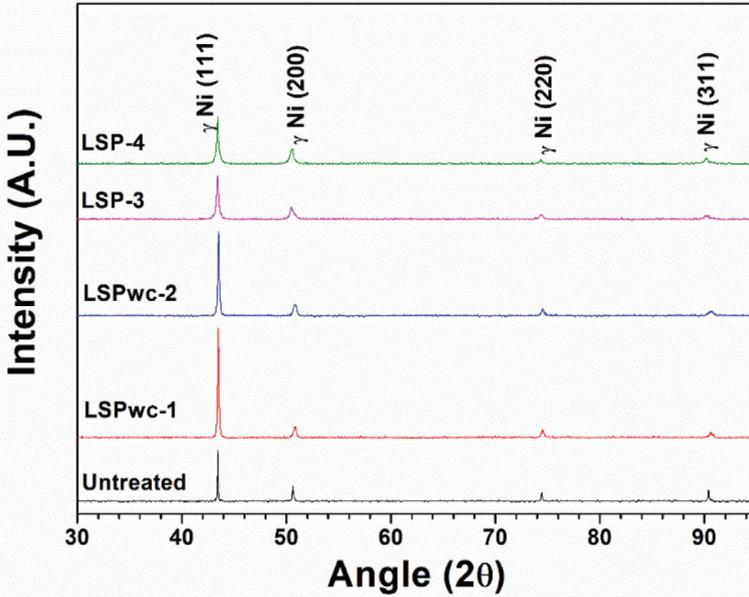


FIGURE 1
XRD scans of untreated and laser shock peened surface under different processing conditions.

loy-X superalloys shows no sign of phase transformation. This is usually expected, and it is indicative that the effects generated herein are similar to cold working. Thus, thermal input as a result of the process was negligible for bringing about such phase transformations which could otherwise be a possibility with thermos/mechanical effect which the LSP process could also produce.

3.2 Estimation of lattice deformation stress and microstrain by the uniform stress deformation model (USDM)

The Williamson-Hall method for uniform deformation model (UDM) is based on the assumption that the crystals are homogeneous and isotropic in nature; that is, the strain is uniform in all the crystallographic directions, and is represented by [19]

$$\beta_{hkl} \cos \theta_{hkl} = \frac{k\lambda}{D} + \epsilon \sin \theta_{hkl} \quad (5)$$

where β is the corrected full width at half maximum (FWHM), θ is the diffraction angle, K (≈ 1) is a constant, λ is the X-ray wavelength and D is the crystallite size (or domain size).

It is unrealistic to assume that the crystals are homogeneous and isotropic in nature or the strain is independent of crystallographic directions. So, we have used uniform stress deformation model (USDm) considering the anisotropic nature of strain. The model assumes a uniform lattice deformation stress. The microstrain in Equation (5) can be rewritten as

$$\varepsilon = \frac{\sigma}{E_{hkl}} \quad (6)$$

where E_{hkl} is the Young's modulus in a crystallographic direction perpendicular to the lattice plain (hkl). The crystallographic dependence of the Young's modulus in a cubic crystal is given by

$$\frac{1}{E_{hkl}} = S_{11} - 2 \left(S_{11} - S_{12} - \frac{1}{2} S_{44} \right) \left[\frac{h^2 k^2 + k^2 l^2 + l^2 h^2}{(h^2 + k^2 + l^2)^2} \right] \quad (7)$$

where S_{11} , S_{12} , and S_{44} are the elastic compliance which, for a Ni alloy, are 7.67, 2.93 and 8.23 TPa⁻¹, respectively [20]. The values of Young's modulus corresponding to different crystallographic planes are given in Tables 3 to 7. Re-writing Equation (5), whilst considering anisotropy in microstrain or lattice deformation stress, gives

$$\beta_{hkl} \cos \theta_{hkl} = \frac{k\lambda}{D} + 4 \left(\frac{\sigma}{E_{hkl}} \right) \sin \theta_{hkl} \quad (8)$$

Figure 2 shows the plots between $\beta_{hkl} \cos \theta$ and $4 \sin \theta / E_{hkl}$ for the untreated Hastelloy-X and the laser peened samples. From the slope and intercept of the plot between $\beta_{hkl} \cos \theta$ and $4 \sin \theta / E_{hkl}$, the lattice deformation stress and the crystallite size have been measured. The measured values of microstrain and lattice deformation stress in an untreated and laser shock peened Hastelloy-X samples are presented in Tables 3 to 7.

The lattice deformation stress in an untreated Hastelloy-X sample is 49 MPa as shown in Table 3. The lattice deformation stress in LSPwc-1 is 290 MPa which is maximum among the other studied systems (see Tables 4 to 7). A high value of lattice deformation stress in a sample implies that the sample has undergone a significant amount of plastic deformation due to LSP which is attributed to the use of higher power density in LSPwc-1 sample. The lattice deformation stress in LSPwc-2 was found to be 283 MPa and is well above the LSP-3 and LSP-4 systems. This was due to a higher portion of the

TABLE 3

Young's modulus, microstrain, and lattice deformation stress in an untreated Hastelloy-X sample.

Plane (hkl)	Young's Modulus (MPa)	Microstrain (ϵ_{hkl})	Lattice Deformation Stress, σ (MPa)
111	229	0.00017	49
200	130	0.00038	
220	226	0.00022	
311	178	0.00028	
331	243	0.00020	

TABLE 4

Young's modulus, microstrain, and lattice deformation stress in LSPwc-1.

Plane (hkl)	Young's Modulus (MPa)	Microstrain (ϵ_{hkl})	Lattice Deformation Stress, σ (MPa)
111	229	0.00097	290
200	130	0.00222	
220	226	0.00128	
311	178	0.00163	
331	243	0.00119	

TABLE 5

Young's modulus, microstrain, and lattice deformation stress in LSPwc-2.

Plane (hkl)	Young's Modulus (MPa)	Microstrain (ϵ_{hkl})	Lattice Deformation Stress, σ (MPa)
111	229	0.00095	283
200	130	0.00217	
220	226	0.00125	
311	178	0.00159	
331	243	0.00116	

laser energy transferred to the sample rather than creating a shock wave with high pressure since the absence of an absorptive layer. The lattice deformation stresses in LSP-3 and LSP-4 systems were measured to be 257 MPa and 198 MPa, respectively. Comparing LSP-3 and LSP-4, it can be concluded that the Al tape as an absorptive layer was effective in creating maximum lattice deformation stress than the black vinyl tape.

TABLE 6
Young's modulus, microstrain, and lattice deformation stress in LSP-3.

Plane (hkl)	Young's Modulus (MPa)	Microstrain (ϵ_{hkl})	Lattice Deformation Stress, σ (MPa)
111	229	0.00086	257
200	130	0.00197	
220	226	0.00114	
311	178	0.00145	
331	243	0.00106	

TABLE 7
Young's modulus, microstrain, and lattice deformation stress in LSP-4.

Plane (hkl)	Young's Modulus (MPa)	Microstrain (ϵ_{hkl})	Lattice Deformation Stress, σ (MPa)
111	229	0.00066	198
200	130	0.00152	
220	226	0.00088	
311	178	0.00112	
331	243	0.00081	

3.3 Dislocation density

3.3.1 Qualitative analysis

Intensity of plastic deformation in the material due to LSP is directly related to the microstrain developed in the crystal lattice. An increase in the microstrain (peak broadening) implies an increase in the dislocation density. LSP of Hastelloy-X superalloys shows that the microstrain (or the plastic deformation) developed on the surface of the Hastelloy-X samples is dependent on the peak power density of the laser as shown in Table 8. With a peak power density of 7.96 GW/cm^2 , the microstrain in the crystal lattice shows a maximum value of 1.460×10^{-3} . Lowest microstrain is measured in the LSP-4 sample where the peak power density of 2.26 GW/cm^2 was used with a black vinyl tape as an absorptive layer. The increase in microstrain in LSPwc-1 is attributed to increase in the shock wave pressure resulting from higher power density according to Equation (1).

3.3.2 Quantitative analysis

Another way to quantify plastic deformation is to measure the dislocation density. Dislocations are the result of plastic deformation in the materials. A higher dislocation density, not only, implies a hardened surface, but also, forms a nanocrystalline structure on the surface which increases its strength.

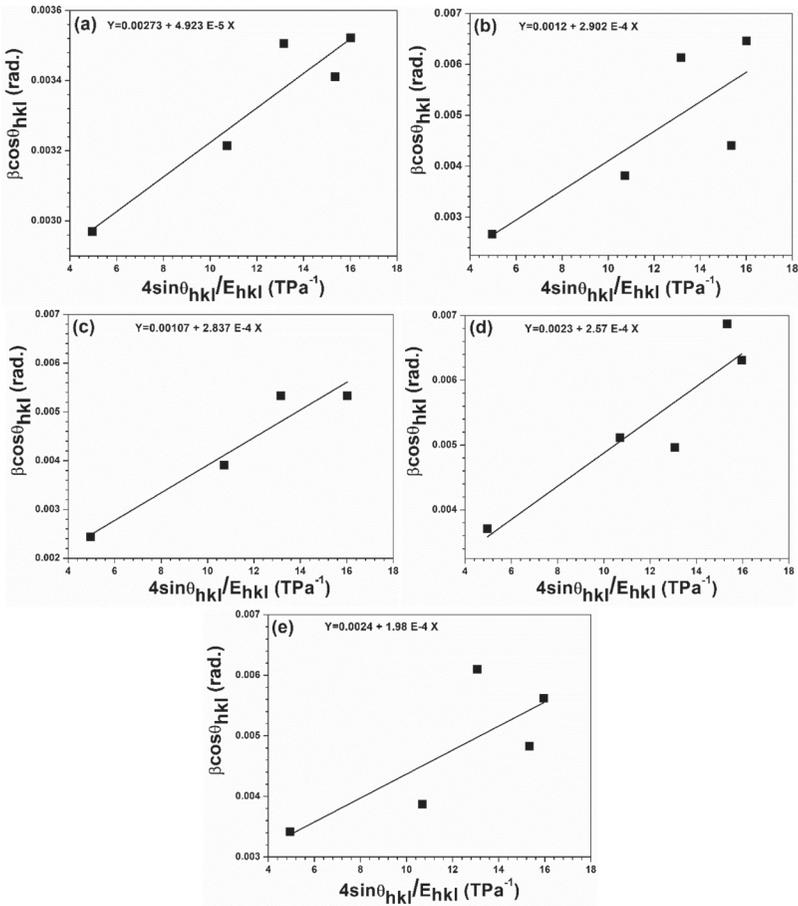


FIGURE 2
Williamson-Hall plots for (a) an untreated Hastelloy-X sample, (b) LSPwc-1, (c) LSPwc-2, (d) LSP-3 and (e) LSP-4.

TABLE 8
Average microstrain and dislocation density in untreated and laser shock peened Hastelloy-X superalloy samples.

Sample	Average Microstrain, $\epsilon \times 10^{-3}$	Dislocation Density, $\rho \times 10^{14} (\text{m}^{-2})$
Untreated	0.250	0.16
LSPwc-1	1.460	5.28
LSPwc-2	1.430	5.06
LSP-3	1.290	4.12
LSP-4	0.998	2.47

In the present study, dislocation density in the untreated and laser shock peened superalloys were measured using Williamson and Smallman approach [18]. This approach considers crystallite size and microstrain as the factors responsible for dislocations density; however, an increase in the dislocation density increases the full width half maximum (FWHM) of the XRD peak which is directly related to the microstrain in the crystal lattice. Table 8 presents the dislocation densities in an untreated and laser shock peened Hastelloy-X superalloys. The LSP of Hastelloy-X superalloys show an increase in the dislocation density as compared to that of the untreated superalloy. This observation indicates that the chosen Nd:YAG laser processing parameters are acceptable and corroborate well with the theory of LSP (*cf.* Equation (1)). Surprisingly, the dislocation density in LSPwc-1 sample is the highest amongst the other laser shock peened samples. The sample laser shock peened with a power density of 7.96 GW/cm^2 shows highest dislocation density.

The use of an absorptive layer in LSP decreases the dislocation density as shown in Table 8. This is not usual as it is believed that the use of absorptive coatings increases the shock wave pressure in the material. The observation can be related to the fact that both thermal and mechanical effects were present in LSPwc-1 and LSPwc-2. The decrease in dislocation densities in LSP-3 and LSP-4 as compared to LSPwc-1 and LSPwc-2 is believed due to the partial loss in the incident laser energy due to melting and evaporation of absorptive coatings in LSP-3 and LSP-4 samples. Comparing LSP-3 and LSP-4, it can be concluded that the Al based absorptive coating increases the dislocation density in the material. This is possibly due to increased shock pressure in the Hastelloy-X superalloys with Al as an absorptive layer as compared to that of black vinyl tape.

3.4 Residual stress analysis

Residual stress developed on the surface of the untreated and laser shock peened samples are presented in Table 9. As shown in Table 9, the residual stress in an untreated Hastelloy-X sample is tensile in nature with 23 MPa and 25 MPa in longitudinal and transverse directions, respectively. Lower values of residual stress measured on the surface of an untreated Hastelloy-X are due to the full annealing operation performed on the samples. The residual stresses developed on the surface of laser shock peened Hastelloy-X samples without absorptive coatings show tensile stresses in both longitudinal and transverse directions despite being the obvious benefits which the process introduces in general.

The reason behind the developed tensile stresses on the surfaces of LSPwc-1 and LSPwc-2 is due to the dominant thermal effect (melting and evaporation) instead of mechanical effect. Normally, an absorptive coating is

TABLE 9
Residual stress developed in the untreated and laser shock peened Hastelloy-X superalloy samples.

Sample	Residual Stress (MPa)	
	σ (Longitudinal)	σ (Transverse)
Untreated	23 ± 13	25 ± 11
LSPwc-1	149 ± 28	135 ± 25
LSPwc-2	121 ± 11	124 ± 22
LSP-3	-166 ± 15	-132 ± 23
LSP-4	-117 ± 21	-48 ± 18

used to avoid thermal effects and ensure a pure mechanical effect develops on the surface of the samples. But, in the case of LSPwc-1 and LSPwc-2, most of the energy was used to melt and evaporate a thin layer of the sample rather than creating shock pressure which is rather needed for inducing compressive residual stress on the surface. Such phenomena was observed by Masse and Barreau [21] and Gill *et al* [22]. On the hand, the residual stresses developed on the surface of the LSP-3 and LSP-4 shows compressive in nature.

The measured compressive nature of the stresses in the LSP-3 and LSP-4 samples are due to the use of the absorptive coatings which prevented the samples surface to get melted and generate the necessary shock pressure in the material in order to induce compressive residual stresses. The developed shock pressure on the surface of the samples are responsible for the formation of compressive residual stresses on the surface of LSP-3 and LSP-4.

It is interesting to note that the surface treated with Al tape shows the maximum residual compressive stress than the surface treated with black vinyl tape as shown in Table 9. The reason behind such a difference in generation of stresses on the surface of the Hastelloy-X samples is related to the developed peak pressure in Al tape as compared to black vinyl tape.

3.5 Microhardness distribution

The improvement in microhardness is related to the strain hardening phenomena due to LSP. Figure 3 shows the microhardness distribution in the laser shock peened samples LSP-3 and LSP-4. The samples LSP-3 and LSP-4 were chosen for the microhardness measurement as they showed compressive residual stress on the surface. From Figure 3, it is evident that the microhardness on the surface is higher than the interior. The tail end of the microhardness curve presents the microhardness value of an untreated superalloy (around 260 VHN). The increased microhardness on the surface of the LSP treated sample is related to the strain hardening of the surface layer due to increased dislocation density (see Table 8).

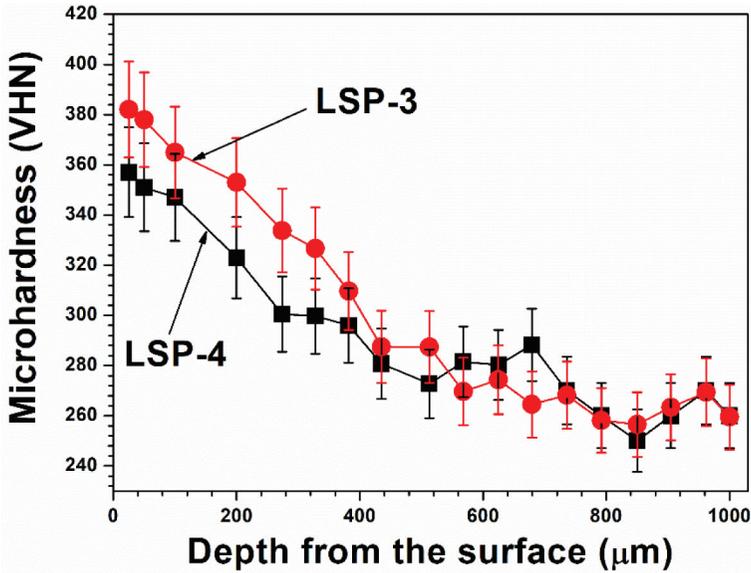


FIGURE 3
Microhardness distribution with depth in LSP-3 (Plot 1) and LSP-4 (Plot 2) Hastelloy-X superalloy sample.

The gradual decrease in the microhardness values of the LSP-3 and LSP-4 samples with depth is related to the decrease in the dislocation density with depth. The material, normally, experiences a maximum shock pressure due to LSP on the surface which gradually decreases towards the depth. As the amount of plastic deformation is dependent on the shock wave pressure, the plastic deformation decreases with depth. The laser shock peened superalloy with Al as an absorptive layer (LSP-3) shows maximum improvement in hardness with around 380 VHN as compared to around 360 VHN in the LSP-4. The improvement in microhardness in LSP-3 is related to the increased dislocation density and higher compressive residual stress.

4 CONCLUSIONS

A comprehensive investigation of laser shock peening (LSP) surface treatment of Hastelloy-X superalloys was undertaken using a 10 J Nd:YAG laser to understand the evolution of structure and properties as well as to understand the strengthening mechanisms. The LSP surface treatment was carried out on the bare Hastelloy-X superalloys as well as on the superalloys coated with Al and black vinyl tape as absorptive layers. The process parameters were chosen to ensure plastic deformation in the material based on the clas-

sical theory of shock wave (2.26 and 7.96 GW/cm²). From the study, the following conclusions may be drawn:

- (i) The LSP treatment showed no sign of phase transformation ensuring that the sample's surface structure was stable after LSP treatment;
- (ii) The lattice deformation stress was found to increase with increase in the peak power density. A higher lattice deformation stress of 290 MPa was obtained for Hastelloy-X samples LSP treated without an absorptive layer;
- (iii) A larger microstrain of 1.520×10^{-3} was developed on the surface of the Hastelloy-X sample with LSP treated at highest power density and without an absorptive tape. Use of Al as an absorptive layer during LSP showed an increase in the microstrain (1.290×10^{-3}) in the crystal lattice of Hastelloy-X superalloy in comparison to the black vinyl tape (0.998×10^{-3});
- (iv) Tensile stresses were developed on the Hastelloy-X superalloys when they were subjected to LSP without any absorptive layer. On the other hand, compressive residual stresses of 48 to 166 MPa were developed on the surface of Hastelloy-X superalloys when they were laser shock peened with aforementioned absorptive layers. Maximum compressive residual stress of 166 MPa was observed on the Hastelloy-X superalloy when laser shock peened with Al absorptive layer;
- (v) Dislocation densities were found to be more when the samples were laser shock peened without an absorptive layers implying higher plastic deformation in the sample. High laser power density was also able to increase the dislocation density in the Hastelloy-X samples; and
- (vi) Strain hardening was more on the samples treated with Al absorptive layer compared to black vinyl tape. A maximum hardness of around 380 VHN was found with Al absorptive layer compared to around 360 VHN for black vinyl tape.

ACKNOWLEDGMENTS

The authors of this paper would like to thank the EPSRC funded laser loan pool scheme for granting a state-of-the-art shock peening laser system (Grant No: 13250017 (NSL4)).

NOMENCLATURE

<i>a</i>	Lattice parameter (nm)
<i>b</i>	Burger's vector (nm)

d_0	Stress-free lattice spacing (nm)
$d\phi_0$	Stress free interplanar lattice spacing (nm)
D	Crystallite size (nm)
E	Young's modulus (GPa)
HEL	Hugoniot elastic limit (GPa)
I_0	Laser power density (GW/cm^2)
P	Shock wave pressure (GPa)
S	Elastic compliance factor (TPa^{-1})
Z	Reduced shock impedance between the material and the confining medium ($\text{g}/\text{cm}^2\text{s}^2$)

Greek symbols

α	Efficiency of the interaction between the material and the confining medium
β	Full width half maximum, FWHM (radian)
ε	Microstrain in a crystal lattice
θ	Diffraction angle ($^\circ$)
λ	X-ray wavelength (A°)
ρ	Dislocation density (m^{-2})
σ_Y^{dyn}	Dynamic yield strength (GPa)
σ	Lattice deformation stress (MPa)
$\sigma\phi$	Residual stress at an angle, ϕ (MPa)
ϕ	Angle between a fixed direction in the plane of the sample and the projection in that plane of the normal of the diffracting plane
Ψ	Angle between the normal of the sample and the normal of the diffracting plane (bisecting the incident and diffracted beams)

REFERENCES

- [1] Peyre P. and Fabbro R. Laser shock processing: A review of the physics and applications. *Optical and Quantum Electronics* **27**(12) (1995), 1213-1229.
- [2] Peyre P., Fabbro R., Merrien P. and Lieurade H.P. Laser shock processing of aluminium alloys. Application to high cycle fatigue behaviour. *Materials Science & Engineering A. Structural Materials: Properties, Microstructure and Processing* **210**(1-2) (1996), 102-113.
- [3] Lu J.Z., Qi H., Luo K.Y., Luo M. and Cheng X.N. Corrosion behaviour of AISI 304 stainless steel subjected to massive laser shock peening impacts with different pulse energies. *Corrosion Science* **80** (2014), 53-59.

- [4] Lim H., Kim P., Jeong H. and Jeong S. Enhancement of abrasion and corrosion resistance of duplex stainless steel by laser shock peening. *Journal of Materials Processing Technology* **212**(6) (2012), 1347–1354.
- [5] Montross C.S., Wei T., Ye L., Clark G. and Mai Y.W. Laser shock processing and its effects on microstructure and properties of metal alloys: A review. *International Journal of Fatigue* **24**(10) (2002), 1021–1036.
- [6] Rubio-González C., Felix-Martinez C., Gomez-Rosas G., Ocana J.L., Morales M. and Porro J.A. Effect of laser shock processing on fatigue crack growth of duplex stainless steel. *Materials Science & Engineering A. Structural Materials: Properties, Microstructure and Processing* **528**(3) (2011), 914–919.
- [7] Shukla P., Nath S., Wang G., Shen X. and Lawrence J. Surface property modifications of silicon carbide ceramic following laser shock peening. *Journal of the European Ceramic Society* **37**(9) (2017), 3027–3038.
- [8] Reed R.C. *The Superalloys: Fundamentals and Applications*. New York: Cambridge University Press. 2006.
- [9] Callister W.D. and Rethwisch D.G. *Materials Science and Engineering*. Hoboken: John Wiley & Sons. 2014.
- [10] James M.N., Hughes D.J., Chen Z., Lombard H., Hattingh D.G., Asquith D., Yates J.R. and Webster P.J. Residual stresses and fatigue performance. *Engineering Failure Analysis* **14**(2) (2007), 384–395.
- [11] Raghavan R., Ayer R., Jin H.W., Marzinsky C.N. and Ramamurthy U. Effect of shot peening on the fatigue life of a Zr-based bulk metallic glass. *Scripta Materialia* **59**(2) (2008), 167–170.
- [12] Shukla P.P. and Lawrence, J. The influence of brightness whilst laser surface processing of a silicon nitride engineering ceramic. *Optics and Lasers in Engineering* **50** (2012), 1746–17512.
- [13] Shukla, P.P. and Lawrence, J. Role of laser beam radiance in different ceramic processing: A two wavelength comparison. *Optics & Laser Technology* **54** (2013), 380–388.
- [14] Shukla, P., Lawrence J. and Zhang, Yu. Understanding laser-beam brightness: a review on a new prospective in materials processing. *Optics and Lasers in Engineering* **75** (2015), 40–51.
- [15] Shukla, P.P. and Lawrence, J. Identification of optical parameters for determination of radiance. *Journal of Optics* **44**(1) (2015), 12–19.
- [16] Fabbro R., Fournier J., Ballard P., Devaux D. and Virmont J. Physical study of laser-produced plasma in confined geometry. *Journal of Applied Physics* **68**(2) (1990), 775–784.
- [17] Johnson J.N. and Rohde R.W. Dynamic deformation twinning in shock-loaded iron. *Journal of Applied Physics* **42**(11) (1971), 4171–4182.
- [18] Williamson G.K. and Smallman R.E. Dislocation densities in some annealed and cold worked metals from measurements on the X-Ray Debye-Scherrer spectrum. *Philosophical Magazine* **1**(1) (1956), 34–46.
- [19] Suryanarayana C. and Norton M.G. *X-Ray Diffraction: A practical approach*. New York: Springer. 1998.
- [20] Gale W.F. and Totemeier T.C. (Eds.). *Smithells Metals Reference Book*. Amsterdam: Elsevier. 2004.
- [21] Masse J.E. and Barreau G. Laser generation of stress waves in metal. *Surface & Coatings Technology* **70**(2-3) (1995), 231–234.
- [22] Gill A.S., Telang, A. and Vasudevan V.K. Characteristics of surface layers formed on Inconel 718 by laser shock peening with and without a protective coating. *Journal of Materials Processing Technology* **225** (2015), 463–472.